

Anatomy of Toroidal Drift Modes

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Toroidal Drift Mode Anatomy (Part I)

Can describe Toroidal ITG, and all three varieties of TEM modes with one dispersion relation.

Ions in Fluid Limit: $\omega_{Di}, k_{\parallel} v_{thi} \ll \omega$

$$\frac{\delta n_i}{n} = -\frac{e\delta\phi}{T_i} \left\{ \frac{\omega_{*i} - \omega_{Di}}{\omega} + \frac{\omega_{*i}(1 + \eta_i)\omega_{Di}}{\omega^2} + \left(b_i - \frac{k_{\parallel}^2 v_{thi}^2}{\omega^2} \right) \left(1 - \frac{\omega_{*i}(1 + \eta_i)}{\omega} \right) \right\}$$

Circulating electrons with $\omega/k_{\parallel} \ll v_{the}$: adiabatic response

$$\frac{\delta n_{eC}}{n_{eC}} = \frac{e\delta\phi}{T_e}$$

Trapped electrons with $v_{the\parallel} \ll \omega/k_{\parallel} \ll v_{the}$

(feel curvature, ExB drifts, but no Landau damping)

$$\partial_t \delta n_{eT} + \nabla \cdot \left(\underbrace{n_{eT} \delta \mathbf{V}_{eT}}_{\text{ExB drift}} + \underbrace{\mathbf{V}_{eT} \delta n_{eT}}_{\text{curvature drift}} \right) = 0$$

$$\frac{\delta n_{eT}}{n_{eT}} = \frac{\omega_{*e}}{\omega - \omega_{De}} \frac{e\delta\phi}{T_e}$$

$$\eta_i = d \ln T_i / d \ln n$$

$$\omega_D = k_{\theta} g / \Omega_c$$

$$g = v_{th}^2 / R$$

$$b_i = k_{\theta}^2 \rho_i^2 / 2$$

$$\omega_* = k_{\theta} D_B / r_n$$

$$D_B = cT / ZeB$$

$$r_n^{-1} = -d \ln n / dr$$

$$\epsilon_n = r_n / R_0$$

$$\tau = T_e / T_i$$

Impose quasineutrality to obtain dispersion relation

$$\delta n_i = \delta n_{eC} + \delta n_{eT}$$

Refs:

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Toroidal Drift Mode Anatomy (Part II)

First, improve on electron response with approximate kinetic theory...

$$\delta f_e \equiv \frac{e\delta\phi}{T_e} F_{me} + \delta h_e$$

$$\left(\omega - \omega_{De} + i\nu_{\text{eff}} + i \frac{v_{\parallel}}{qR} \frac{\partial}{\partial \theta} \right) \delta h_e = - \frac{e\delta\phi}{T_e} F_{me} (\omega - \omega_{*e})$$

$$\oint \frac{d\theta}{v_{\parallel}} (\dots) \Rightarrow \delta h_e \simeq - \frac{e\delta\phi}{T_e} F_{me} \frac{\omega - \omega_{*e} (1 + \eta_e (E/T_e - 3/2))}{\omega - \omega_{De} + i\nu_{\text{eff}}}$$

Integral over velocity space gives perturbed density:

$$\frac{\delta n_{eC}}{n_{eC}} = \frac{e\delta\phi}{T_e} \quad \langle \dots \rangle = \frac{2}{\sqrt{\pi}} \int_0^{\infty} dx \sqrt{x} e^{-x} (\dots)$$

$$\frac{\delta n_{eT}}{n_{eT}} = \frac{e\delta\phi}{T_e} \left(1 - \left\langle \frac{\omega - \omega_{*e} (1 + \eta_e (E/T_e - 3/2))}{\omega - \omega_{De} + i\nu_{\text{eff}}} \right\rangle_E \right)$$

$$\frac{\delta n_e}{n_e} = \frac{e\delta\phi}{T_e} \left(1 - \frac{n_{eT}}{n_e} \left\langle \frac{\omega - \omega_{*e} (1 + \eta_e (E/T_e - 3/2))}{\omega - \omega_{De} + i\nu_{\text{eff}}} \right\rangle_E \right)$$

Toroidal Drift Mode Anatomy: Resonant Modes

Construct Dispersion Relation:

$$\frac{\omega_{*i} - \omega_{Di}}{\omega} + \frac{\omega_{*i}(1 + \eta_i)\omega_{Di}}{\omega^2} + \left(b_i - \frac{k_{\parallel}^2 v_{thi}^2}{\omega^2}\right) \left(1 - \frac{\omega_{*i}(1 + \eta_i)}{\omega}\right) + \frac{1}{\tau} - \frac{1}{\tau} \frac{n_{eT}}{n} \left\langle \frac{\omega - \omega_{*e}(1 + \eta_e(E/T_e - 3/2))}{\omega - \omega_{De} + i\nu_{eff}} \right\rangle_E = 0$$

1) Dissipative Trapped Electron Mode - destabilized by collisions

$$\frac{\omega_{*i} - \cancel{\omega_{Di}}}{\omega} + \frac{\omega_{*i}(1 + \eta_i)\cancel{\omega_{Di}}}{\cancel{\omega^2}} + \left(b_i - \frac{k_{\parallel}^2 v_{thi}^2}{\cancel{\omega^2}}\right) \left(1 - \frac{\omega_{*i}(1 + \eta_i)}{\cancel{\omega}}\right) + \frac{1}{\tau} - \frac{1}{\tau} \frac{n_{eT}}{n} \left\langle \frac{\omega - \omega_{*e}(1 + \eta_e(E/T_e - 3/2))}{\omega - \cancel{\omega_{De}} + i\nu_{eff}} \right\rangle_E = 0$$

$$\gamma_{DTEM} \propto \frac{\epsilon^{3/2} \omega_{*e}^2}{\nu_{ei}} \left(\frac{3}{2} \eta_e + b_s\right)$$

2) Collisionless Trapped Electron Mode - destabilized by resonant denominator

$$\frac{\omega_{*i} - \cancel{\omega_{Di}}}{\omega} + \frac{\omega_{*i}(1 + \eta_i)\cancel{\omega_{Di}}}{\cancel{\omega^2}} + \left(b_i - \frac{k_{\parallel}^2 v_{thi}^2}{\cancel{\omega^2}}\right) \left(1 - \frac{\omega_{*i}(1 + \eta_i)}{\cancel{\omega}}\right) + \frac{1}{\tau} - \frac{1}{\tau} \frac{n_{eT}}{n} \left\langle \frac{\omega - \omega_{*e}(1 + \cancel{\eta_e}(E/T_e - 3/2))}{\omega - \omega_{De} + i\cancel{\nu_{eff}}}\right\rangle_E = 0$$

$$\gamma_{CTEM} \simeq \omega_{*e} \left(\frac{2\pi r}{R}\right)^{1/2} \eta_e \left(\frac{R}{r_n}\right)^{3/2} \left(\frac{R}{r_n} - \frac{3}{2}\right) \exp(-R/r_n)$$

Toroidal Drift Mode Anatomy: Non-Resonant "Fluid" Modes

Expand electron response for $\omega_{De} \ll \omega \Rightarrow \frac{\omega - \omega_{*T}}{\omega - \omega_D} \simeq \left(1 - \frac{\omega_{*T}}{\omega}\right) \left(1 + \frac{\omega_D}{\omega}\right)$

$$\frac{\delta n_e}{n} \simeq \frac{e\delta\phi}{T_e} \left(1 - \frac{n_{eT}}{n} \left[1 - \frac{\omega_{*e}(1 + \eta_e)}{\omega} + \frac{\omega_{De}}{\omega} - \frac{\omega_{*e}(1 + \eta_e)\omega_{De}}{\omega^2}\right]\right)$$

Fluid Dispersion Relation:

$$\begin{aligned} & \left(1 - \frac{n_{eT}}{n} + \tau b_i\right) \\ & + \frac{\omega_{*e}}{\omega} \left(\frac{n_{eT}}{n}(1 + \eta_e) + \frac{r_n}{R}(1 - n_{eT}/n) + b_i(1 + \eta_i) - 1\right) \Rightarrow \\ & + \frac{\omega_{*e}^2}{\omega^2} \left(\frac{n_{eT} r_n}{n R}(1 + \eta_e) + b_i(1 + \eta_i) + \frac{1 + \eta_i r_n}{\tau R}\right) = 0 \end{aligned}$$

Choose b_i to eliminate this term, leaving purely growing modes - gives wavelength for strongest modes (notice real frequency depends on b_i and η_i)

For small, η_i fluid ion response not strictly valid...but...

$$\gamma \simeq \omega_{*e} \sqrt{\frac{\frac{n_{eT} r_n}{n R}(1 + \eta_e) + b_i(1 + \eta_i) + \frac{1 + \eta_i r_n}{\tau R}}{1 - n_{eT}/n + \tau b_i}}$$

$$\gamma_{TEM} \sim (k_\theta \rho_i) \sqrt{(n_T/n) g_i / r_n}$$

"Ubiquitous mode"

$$\gamma_{ITG} \sim (k_\theta \rho_i) \sqrt{\eta_i g_i / r_n}$$

Rayleigh-Taylor analog: $g = v_{th}^2 / R$